

## Magnetic properties and magnetocaloric effects of $RNiSi_2$ (R= Gd, Dy, Ho, Er, Tm) compounds

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## Magnetic properties and magnetocaloric effects of RNiSi<sub>2</sub> (R= Gd, Dy, Ho, Er, Tm) compounds

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Orthorhombic CeNiSi<sub>2</sub>-type polycrystalline RNiSi<sub>2</sub> (R=Gd, Dy, Ho, Er, Tm) compounds were synthesized and the magnetic and magnetocaloric properties were investigated in detail. The transition temperatures of RNiSi<sub>2</sub> compounds are all in a very low temperature range (<30 K). As temperature increases, all of the compounds undergo an AFM to PM transition (GdNiSi<sub>2</sub> at 18 K, DyNiSi<sub>2</sub> at 25 K, HoNiSi<sub>2</sub> at 10.5 K, ErNiSi<sub>2</sub> at 3 K and TmNiSi<sub>2</sub> at 3.5 K, respectively). ErNiSi<sub>2</sub> compound shows the largest  $(\Delta S_M)_{max}$  (maximal magnetic entropy change) among these compounds. The value of  $(\Delta S_M)_{max}$  is 27.9 J/kgK under a field change of 0-5 T, which indicates that ErNiSi<sub>2</sub> compound is very competitive for practical applications in low-temperature magnetic refrigeration in the future. DyNiSi<sub>2</sub> compound shows large inverse MCE (almost equals to the normal MCE) below the T<sub>N</sub> which results from metamagnetic transition under magnetic field. Considering of the normal and inverse MCE, DyNiSi<sub>2</sub> compound also has potential applications in low-temperature multistage refrigeration. © 2018 Author(s). All article content, except where otherwise noted, is licensed under a Creative Commons Attribution (CC BY) license (<http://creativecommons.org/licenses/by/4.0/>). <https://doi.org/10.1063/1.5007018>

### I. INTRODUCTION

The magnetocaloric effect (MCE), a kind of physical phenomenon found more than 130 years ago, is one of the intrinsic properties of magnetic materials. Commercial and residential refrigeration, based on conventional gas compression/expansion technology, is a mature industry. However, the magnetic refrigeration (MR) based on MCE is becoming competitive with conventional gas compression/expansion technology because its environmental friendliness and high efficiency.<sup>1-5</sup> In the last twenty years great progress has been made on exploring large MCE materials for applications at room temperature such as refrigerators and air conditioners.<sup>6,7</sup> The typical room temperature MCE materials mainly include Gd<sub>5</sub>Si<sub>2</sub>Ge<sub>2</sub>,<sup>8,9</sup> La(Fe, Si)<sub>13</sub>,<sup>10-14</sup> MnAs<sub>1-x</sub>Sb<sub>x</sub>,<sup>15</sup> MnFeP<sub>1-x</sub>As<sub>x</sub>,<sup>16</sup> Heusler alloys,<sup>17,18</sup> etc. MCE materials are generally evaluated by the following parameters: the maximum value of magnetic entropy change  $(\Delta S_M)_{max}$ , the full width at half maximum of  $\Delta S_M - T$  curve ( $\delta T_{FWHM}$ ), refrigerant capacity (RC) and adiabatic temperature change ( $\Delta T_{ad}$ ). For most MCE materials, the value of the  $\Delta S_M$  is negative (normal MCE). But the positive  $\Delta S_M$  which called inverse MCE also can be found in some materials such as TbNiGe<sub>2</sub>,<sup>19</sup> TbMn<sub>0.33</sub>Ge<sub>2</sub>,<sup>20</sup> etc. Recently, much attention has also been paid on the MCE materials with low transition temperature because these materials are promising to be used for gas liquefaction in magnetic cooling cycle or combined magnetic-gas cooling cycle.<sup>6,21</sup> These low temperature MCE materials mainly include the rare earth based intermetallic compounds such as RCo<sub>2</sub>,<sup>22,23</sup> RNi,<sup>24</sup> etc.

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Rare earth based intermetallic compounds have shown interesting magnetic properties and excellent performance on MCE. The ternary  $\text{RTSi}_2$  series (R is rare earth, T is transition metal) is one category of rare earth based intermetallic compounds which has been intensively studied during the last decades.<sup>25–27</sup> The results show that the compounds with  $\text{R}=\text{Pr}$  and  $\text{Nd}$  are ferromagnets and those with  $\text{R}=\text{Gd}$ ,  $\text{Tb}$ ,  $\text{Dy}$ ,  $\text{Ho}$ ,  $\text{Er}$  are antiferromagnets. Experimental results of electrical resistivity, Hall coefficient, magnetic susceptibility and specific heat for  $\text{CeNiSi}_2$  support a theory of the Anderson lattice.<sup>28</sup> Neutron diffraction and magnetic measurements on  $\text{RNiSi}_2$  compounds where  $\text{R}=\text{Pr}$  and  $\text{Nd}$  give refined atomic position parameters and it is found that the preferred moment direction is along the  $c$  axis.<sup>29</sup> The same measurements on  $\text{RNiSi}_2$  compounds where  $\text{R}=\text{Tb}$ ,  $\text{Dy}$ ,  $\text{Ho}$ ,  $\text{Er}$  show that the compounds all have uniaxial moment arrangements and have almost the same easy axis, the shortest axis  $c$ , except for  $\text{ErNiSi}_2$  with an easy axis of a axis.<sup>30–33</sup>

Considering the lack of research work on MCE of  $\text{RNiSi}_2$  compounds, further study on magnetic properties and MCE will be performed. In this paper, the polycrystalline  $\text{RNiSi}_2$  ( $\text{R}=\text{Gd}$ ,  $\text{Dy}$ ,  $\text{Ho}$ ,  $\text{Er}$ ,  $\text{Tm}$ ) compounds were synthesized. The magnetic properties and MCE which both contains normal and inverse effects were investigated in detail.

## II. EXPERIMENTAL PROCEDURE

The polycrystalline  $\text{RNiSi}_2$  ( $\text{R}=\text{Gd}$ ,  $\text{Dy}$ ,  $\text{Ho}$ ,  $\text{Er}$ ,  $\text{Tm}$ ) compounds were prepared by arc-melting appropriate proportion of constituent components with the purity better than 99.9% in a water-cooled copper hearth in a high-purity argon atmosphere. The ingots were turned over after each melting and re-melted several times to ensure the compositional homogeneity. After arc-melting, the ingots were wrapped by molybdenum foil respectively and sealed in a high-vacuum quartz tube, annealed at 1273K for 25 days, finally quenched into liquid nitrogen. The crystal structure was characterized by powder X-ray diffraction (XRD) method with  $\text{Cu K}\alpha$  radiation. Magnetic measurements including the temperature dependence of magnetization ( $M$ - $T$ ) and the field dependence of magnetization ( $M$ - $H$ ) curves were performed by employing Vibrating Sample Magnetometer with Quantum Design (SQUID-VSM). Heat capacity measurements were carried out by employing Physical Properties Measurement System (PPMS).

## III. RESULTS AND DISCUSSION

The XRD pattern of  $\text{RNiSi}_2$  ( $\text{R}=\text{Gd}$ ,  $\text{Dy}$ ,  $\text{Ho}$ ,  $\text{Er}$ ,  $\text{Tm}$ ) compounds at room temperature and its crystal structure are shown in Fig. 1. Almost all of the diffraction peaks can be indexed to an

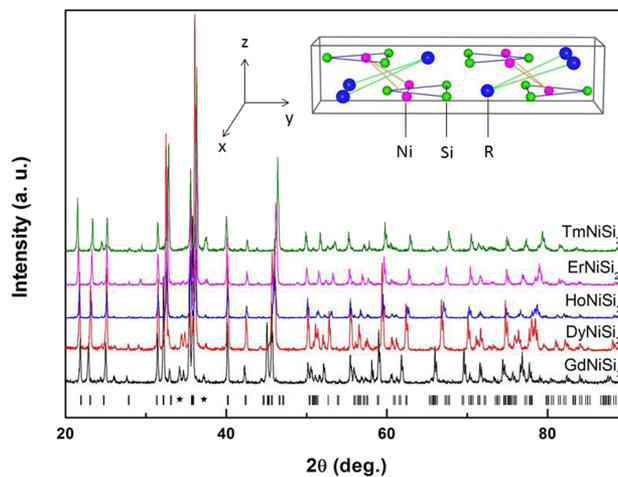


FIG. 1. The XRD patterns of  $\text{RNiSi}_2$  ( $\text{R}=\text{Gd}$ ,  $\text{Dy}$ ,  $\text{Ho}$ ,  $\text{Er}$ ,  $\text{Tm}$ ) compounds measured at room temperature. The inset is the crystal structure of  $\text{RNiSi}_2$ .

orthorhombic CeNiSi<sub>2</sub>-type structure (space group Cmc#63). The result is in accord with previous work.<sup>26</sup> The Bragg positions are marked at the bottom of the picture. It can also be seen that there is a small peak around 37.5° for all compounds and a peak around 34.5° for GdNiSi<sub>2</sub> and DyNiSi<sub>2</sub> compound, which indicates that small amount of impurity may exist. The impurity is indexed as RNiSi<sub>3</sub>. However, it does not affect our discussions and conclusions because the amount of impurity is not large. As the atom number of R increases, the position of diffraction peaks moves towards higher angle range, which indicates that the lattice constant becomes smaller from GdNiSi<sub>2</sub> to TmNiSi<sub>2</sub> compound.

The Zero-Field-Cooled (ZFC) and Field-Cooled (FC) magnetization curves for RNiSi<sub>2</sub> compounds were measured under a field of 0.01 T. They are shown in Fig. 2(a)–(d). The MH curves for RNiSi<sub>2</sub> compounds are also shown in Fig. 2(e)–(h). The MT for DyNiSi<sub>2</sub> shows a rapid increase and then a decrease with the increasing temperature, which indicates that this compound undergoes a simple transition from antiferromagnetic (AFM) to paramagnetic (PM) phase. The same conclusion can also be found in MH curves. The transition temperature is determined to be T<sub>N</sub>=25 K. The overlap of ZFC and FC curves around T<sub>N</sub> shows a good thermal reversibility in this compound. From MH curve we can see the compound occurs metamagnetic transition with the field increasing. All of the features mentioned above can also be found in other compounds. The transition temperatures along with the effective magnetic moments, which have been calculated according to the Curie-Weiss Law, are shown in Table I. The effective moment and ion moment is almost the same for each compound, indicating only rare earth atoms contribute to the magnetic moments in this series of compounds. It reveals from Fig. 2 that the value of the transition temperature (T<sub>N</sub>) shows a decreasing trend when the atomic number of rare earth atom increases. However, there is an exception for TmNiSi<sub>2</sub>, which may result from the complex magnetic coupling in TmNiSi<sub>2</sub> compound.

The MCE materials are generally evaluated by isothermal magnetic entropy change ( $\Delta S_M$ ), which is calculated from isothermal magnetization data (M-H curves) by using Maxwell relation:  $\Delta S_M = \int_0^H (\partial M / \partial T)_H dH$ . The temperature dependences of  $\Delta S_M$  under a field change of 0-2 T and 0-5 T for RNiSi<sub>2</sub> compounds are shown in Fig. 3(a)–(d), respectively. The results shows that all compounds show normal MCE because the value of  $\Delta S_M$  is negative around transition temperatures

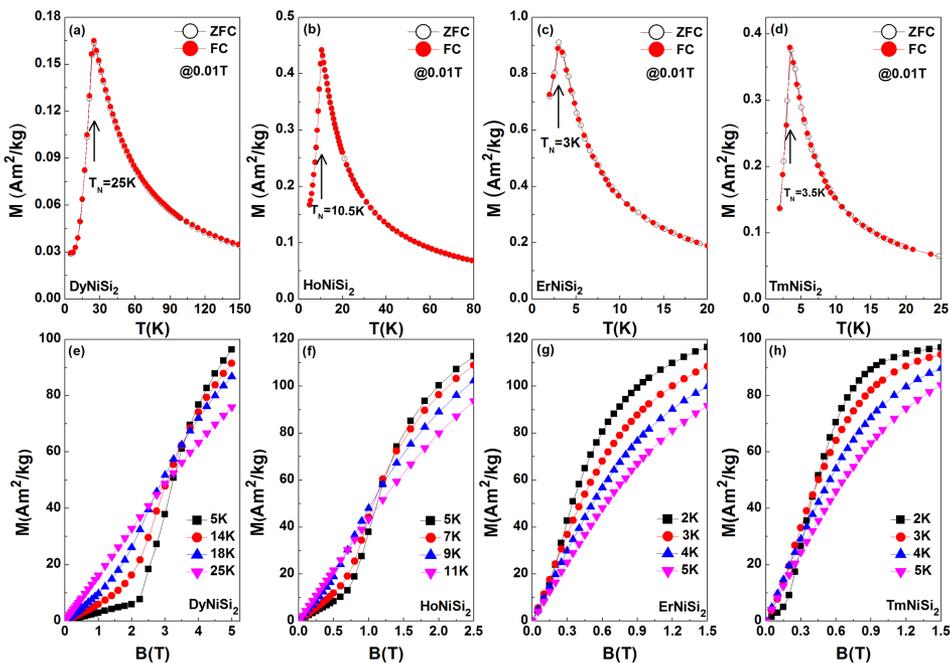
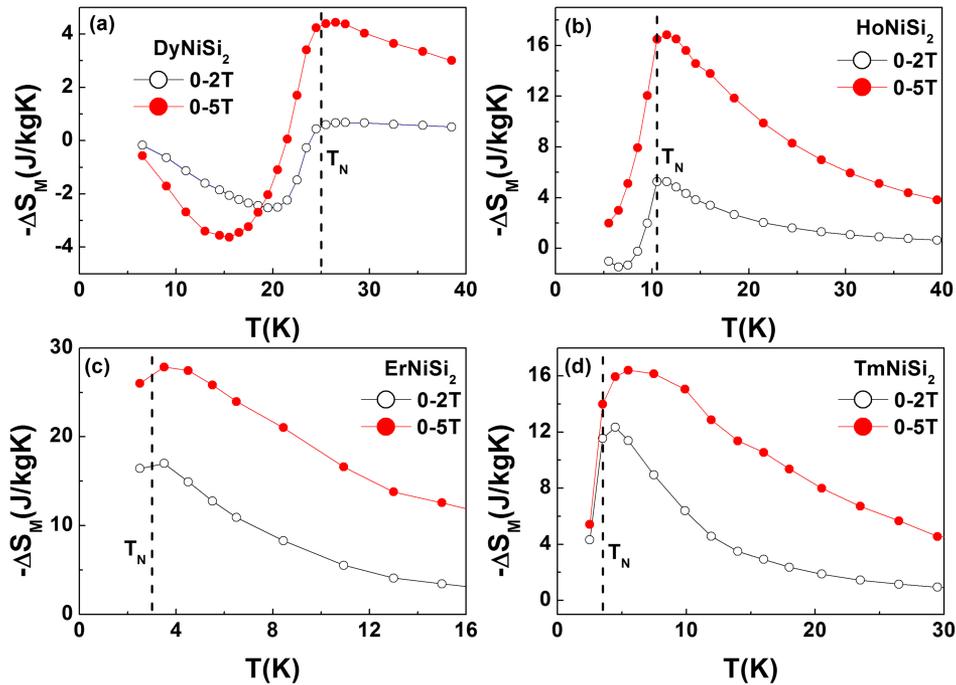


FIG. 2. The MT and MH curves of RNiSi<sub>2</sub> (R= Dy, Ho, Er, Tm) compounds, respectively. (a) MT for DyNiSi<sub>2</sub>. (b) MT for HoNiSi<sub>2</sub>. (c) MT for ErNiSi<sub>2</sub>. (d) MT for TmNiSi<sub>2</sub>. (e) MH for DyNiSi<sub>2</sub>. (f) MH for HoNiSi<sub>2</sub>. (g) MH for ErNiSi<sub>2</sub>. (h) MH for TmNiSi<sub>2</sub>.

TABLE I. The transition temperatures, effective magnetic moments, ion magnetic moments and magnetocaloric parameters of RNiSi<sub>2</sub> (R= Gd, Dy, Ho, Er, Tm) compounds and other compounds.

Materials	$T_N(K)$	$M_{eff}(\mu_B)$	$M_{ion}(\mu_B)$	0-2T			0-5T		
				$-(\Delta S_M)_{max}$ (J/kg K)	$\delta T_{FWHM}(K)$	$RC(J/kg)$	$-(\Delta S_M)_{max}$ (J/kg K)	$\delta T_{FWHM}(K)$	$RC(J/kg)$
GdNiSi <sub>2</sub>	18.0	8.2	7.8	0.7	20.2	9.9	3.5	21.2	58.0
DyNiSi <sub>2</sub>	25.0	10.7	10.5	0.7	15.3	8.6	4.4	15.8	59.0
DyNiSi <sub>2</sub>	-	-	-	(-2.5)	-	31.6	(-3.6)	-	90.5
HoNiSi <sub>2</sub>	10.5	11	10.6	5.3	8.9	35.4	16.8	15.6	198.8
ErNiSi <sub>2</sub>	3.0	9.1	9.6	17.0	-	75.8	27.9	-	225.3
TmNiSi <sub>2</sub>	3.5	7.0	7.1	12.3	7.5	71.2	16.4	17.2	222.6
ErMn <sub>2</sub> Si <sub>2</sub> <sup>34</sup>	4.5	-	-	20.0	-	130	25.2	-	365
TmCuAl <sup>35</sup>	2.8	-	-	17.2	-	129	24.3	-	372
ErNi <sub>2</sub> Si <sub>2</sub> <sup>36</sup>	3.5	-	-	15.1	-	-	22.9	-	-
ErNiSi <sup>37</sup>	3.2	-	-	8.8	-	-	19.1	-	309

for each of them. However, HoNiSi<sub>2</sub> shows small inverse MCE below the transition temperature and DyNiSi<sub>2</sub> shows large inverse MCE (almost equals to the normal MCE) below the transition temperature which will be discussed in detail in the following section. It can be observed from Fig. 3 that all of the curves have a large peak and the maximal magnetic entropy change  $(\Delta S_M)_{max}$  occurs around  $T_N$ . The refrigerant capacity (RC) is another important parameter to evaluate MCE materials. The value of RC can be calculated by using the approach  $RC = \int_{T_1}^{T_2} |\Delta S_M| dT$ , where  $T_1$  and  $T_2$  are the temperatures corresponding to the full width at the half value of  $(\Delta S_M)_{max}$ , respectively. And we call  $\delta T_{FWHM} = T_2 - T_1$  the refrigerant temperature width. The  $(\Delta S_M)_{max}$ , RC and  $\delta T_{FWHM}$  of RNiSi<sub>2</sub> compounds under a field change of 0-2 T and 0-5 T are calculated and shown in Table I. The MCE parameters of some low temperature MCE materials are also shown in Table I for comparison.<sup>34-37</sup> For ErNiSi<sub>2</sub> compound, the value of  $(\Delta S_M)_{max}$  is 17.0 J/kgK for 0-2 T and 27.9 J/kgK for 0-5 T

FIG. 3. The temperature dependences of  $\Delta S_M$  at a field change of 0-2 T and 0-5 T for RNiSi<sub>2</sub> (R= Dy, Ho, Er, Tm) compounds, respectively. (a) DyNiSi<sub>2</sub>. (b) HoNiSi<sub>2</sub>. (c) ErNiSi<sub>2</sub>. (d) TmNiSi<sub>2</sub>.

respectively, which are comparable or even larger than those of other listed materials. The RC of  $\text{ErNiSi}_2$  compound is approximately calculated to be 75.8 J/kg for 0-2 T and 225.3 J/kg for 0-5 T, respectively, where the integration starts with the temperature of 2 K. These values are not accurate because the actual low temperature boundary of  $\delta T_{\text{FWHM}}$  is far lower than 2 K for  $\text{ErNiSi}_2$  compound. According to the model proposed by Oesterrreicher et al.<sup>38</sup> The  $(\Delta S_M)_{\text{max}}$  is positively correlated with the total angular momentum quantum number (J) and negatively correlated with magnetic ordering temperature. In this series,  $\text{HoNiSi}_2$  compound shows the largest J, but its Neel temperature is much larger than that of  $\text{ErNiSi}_2$  compound. Considering that  $\text{HoNiSi}_2$  and  $\text{ErNiSi}_2$  have a similar J, the value of Neel temperature exerts a main effect on the value of  $(\Delta S_M)_{\text{max}}$ . As a result,  $\text{ErNiSi}_2$  compound shows the largest  $(\Delta S_M)_{\text{max}}$  among  $\text{RNiSi}_2$  compounds. The excellent MCE performance of  $\text{ErNiSi}_2$  compound indicates its potential applications in low temperature refrigeration.

The Fig. 3 also shows that  $\text{DyNiSi}_2$  and  $\text{HoNiSi}_2$  compounds both have positive  $\Delta S_M$  below  $T_N$ , which called inverse MCE. This phenomenon results from the mixed exchange interaction and the applied magnetic field leads to a further spin-disordered state,<sup>39</sup> which occurs metamagnetic transition. This makes  $\partial M/\partial T$  positive under  $T_N$ . For  $\text{HoNiSi}_2$  compound, the positive  $\Delta S_M$  can only be found under a low field change and the value of positive  $\Delta S_M$  is far smaller than the absolute value of negative  $\Delta S_M$  around  $T_N$ . When the field change is 5T, the positive  $\Delta S_M$  has been disappeared, which indicates that the AFM ground state of  $\text{HoNiSi}_2$  compound below  $T_N$  is relatively weak. The same phenomenon can also be found in many Heusler alloys and AFM magnetocaloric effect materials.<sup>39</sup> However, for  $\text{DyNiSi}_2$  compound, the inverse MCE can even be found in a high field change such as 5T and the maximal positive  $\Delta S_M$  is almost equals to the maximal negative  $\Delta S_M$ . For further study, heat capacity was measured and  $\Delta S_M$  also calculated using heat capacity data through the expression  $\Delta S_M = \int_0^H \{[C(T, H) - C(T, 0)]/T\} dT$ . The results all shown in Fig. 4. The Fig. 4 shows that the curves obtained using magnetization data are in good agreement with the corresponding curves obtained using heat capacity. And the field change of the maximal positive  $\Delta S_M$  occurs is about 3.6T. The large inverse MCE results from the strong AFM coupling in  $\text{DyNiSi}_2$  compounds which may be related to the large magnetocrystalline anisotropy of Dy atoms. The inverse MCE can also be used for magnetic refrigeration if only the refrigerator works in a reverse process. For example, for  $\text{DyNiSi}_2$  compound, magnetic refrigeration can be realized in a usual working recycle around 25K and it can be used in a reverse working recycle around 15K, which seems useful in multistage refrigeration. And the RC of  $\text{DyNiSi}_2$  compound is re-calculated by considering of the contribution of inverse MCE at lower temperatures which also shown in Table I. The value of RC for a field change of 0-5 T is modified from 59 J/kg and 90.5 J/kg.

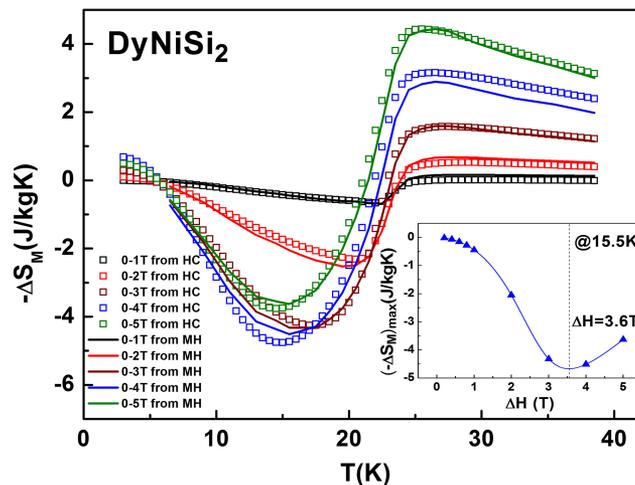


FIG. 4. The temperature dependences of  $\Delta S_M$  for  $\text{DyNiSi}_2$  compound calculated from magnetization and heat capacity data for field changes of 0-1T, 0-2T, 0-3T, 0-4T, and 0-5T, respectively. The inset shows the field change dependences of  $(\Delta S_M)_{\text{max}}$  for inverse MCE at 15.5K.

The excellent performance of DyNiSi<sub>2</sub> compound indicates its potential applications in low temperature refrigeration.

#### IV. CONCLUSION

In summary, all orthorhombic CeNiSi<sub>2</sub>-type polycrystalline RNiSi<sub>2</sub> (R=Gd, Dy, Ho, Er, Tm) compounds are AFM ordered and T<sub>N</sub> in a very low temperature. Among these compounds, ErNiSi<sub>2</sub> compound shows the largest  $(\Delta S_M)_{max}$  and RC, which is larger than almost all of the materials in this low temperature range, indicating its potential practical applications in low-temperature magnetic refrigeration in the future. DyNiSi<sub>2</sub> compound shows large inverse MCE (almost equals to the normal MCE) below the T<sub>N</sub> which results from metamagnetic transition under magnetic field. Considering of the normal and inverse MCE, DyNiSi<sub>2</sub> compound also has potential applications in low-temperature multistage refrigeration.

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- <sup>1</sup> E. Brück, *J. Phys. D: Appl. Phys.* **38**, R381 (2005).
- <sup>2</sup> K. A. Gschneidner and V. K. Pecharsky, *Annu. Rev. Mater. Sci.* **30**, 387 (2000).
- <sup>3</sup> X. Moya, S. Kar-Narayan, and N. D. Mathur, *Nat. Mater.* **13**, 439 (2014).
- <sup>4</sup> V. K. Pecharsky and K. A. Gschneidner, Jr., *J. Magn. Magn. Mater.* **200**, 44 (1999).
- <sup>5</sup> A. M. Tishin and Y. I. Spichkin, *The Magnetocaloric Effect and Its Applications* (2003).
- <sup>6</sup> K. A. Gschneidner, Jr., V. K. Pecharsky, and A. O. Tsokol, *Reports Prog. Phys.* **68**, 1479 (2005).
- <sup>7</sup> B. G. Shen, J. R. Sun, F. X. Hu, H. W. Zhang, and Z. H. Cheng, *Adv. Mater.* **21**, 4545 (2009).
- <sup>8</sup> V. K. Pecharsky and K. A. Gschneidner, Jr., *Phys. Rev. Lett.* **78**, 4494 (1997).
- <sup>9</sup> A. Giguère, M. Foldeaki, B. Ravi Gopal, R. Chahine, T. K. Bose, A. Frydman, and J. A. Barclay, *Phys. Rev. Lett.* **83**, 2262 (1999).
- <sup>10</sup> A. Fujita, S. Fujieda, Y. Hasegawa, and K. Fukamichi, *Phys. Rev. B* **67**, 104416 (2003).
- <sup>11</sup> A. Fujita, Y. Akamatsu, and K. Fukamichi, *J. Appl. Phys.* **85**, 4756 (1999).
- <sup>12</sup> F. X. Hu, B. G. Shen, and J. R. Sun, *Appl. Phys. Lett.* **80**, 826 (2002).
- <sup>13</sup> F. X. Hu, B. G. Shen, J. R. Sun, Z. H. Cheng, G. H. Rao, and X. X. Zhang, *Appl. Phys. Lett.* **78**, 3675 (2001).
- <sup>14</sup> B.-G. Shen, F.-X. Hu, Q.-Y. Dong, and J.-R. Sun, *Chinese Phys. B* **22**, 17502 (2013).
- <sup>15</sup> H. Wada and Y. Tanabe, *Appl. Phys. Lett.* **79**, 3302 (2001).
- <sup>16</sup> O. Tegus, E. Brück, K. H. J. Buschow, and F. R. de Boer, *Nature* **415**, 150 (2002).
- <sup>17</sup> F. Hu, B. Shen, J. Sun, and G. Wu, *Phys. Rev. B* **64**, 132412 (2001).
- <sup>18</sup> J. Liu, T. Gottschall, K. P. Skokov, J. D. Moore, and O. Gutfleisch, *Nat. Mater.* **11**, 620 (2012).
- <sup>19</sup> S. Gupta, K. G. Suresh, A. V. Lukoyanov, Yu. V. Knyazev, and Yu. I. Kuz'min, *J. Alloys Compd.* **664**, 120 (2016).
- <sup>20</sup> S. Gupta and K. G. Suresh, *Physica B* **448**, 260 (2014).
- <sup>21</sup> H. Yayama, Y. Hata, Y. Makimoto, and A. Tomokiyo, *Jpn. J. Appl. Phys.* **39**, 4220 (2000).
- <sup>22</sup> N. H. Duc and D. T. K. Anh, *J. Magn. Magn. Mater.* **242**, 873 (2002).
- <sup>23</sup> A. Giguere, M. Foldeaki, W. Schnelle, and E. Gmelin, *J. Phys. Condens. Matter* **11**, 6969 (1999).
- <sup>24</sup> X. Q. Zheng, B. Zhang, H. Wu, F. X. Hu, Q. Z. Huang, and B. G. Shen, *J. Appl. Phys.* **120**, 163907 (2016).
- <sup>25</sup> A. Gil, A. Szytula, Z. Tomkowicz, K. Wojciechowski, and A. Zygmunt, *Acta Phys. Pol. A* **85**, 271 (1994).
- <sup>26</sup> A. Gil, A. Szytula, Z. Tomkowicz, K. Wojciechowski, and A. Zygmunt, *J. Magn. Magn. Mater.* **129**, 271 (1994).
- <sup>27</sup> V. Ivanov, L. Vinokurova, and A. Szytula, *J. Alloys Compd.* **218**, L24 (1995).
- <sup>28</sup> E. D. Mun, Y. S. Kwon, and M. H. Jung, *Phys. Rev. B* **67**, 33103 (2003).
- <sup>29</sup> P. Schobinger-Papamantellos and K. H. J. Buschow, *J. Alloys Compd.* **185**, 51 (1992).
- <sup>30</sup> P. Schobinger-Papamantellos and K. H. J. Buschow, *J. Less Common Met.* **171**, 321 (1991).
- <sup>31</sup> P. Schobinger-Papamantellos, C. Ritter, and K. H. Buschow, *J. Alloys Compd.* **264**, 89 (1998).
- <sup>32</sup> P. Schobinger-Papamantellos, F. Fauth, and K. H. J. Buschow, *J. Alloys Compd.* **252**, 50 (1997).
- <sup>33</sup> P. Schobinger-Papamantellos, K. H. J. Buschow, C. Wilkinson, F. Fauth, and C. Ritter, *J. Magn. Magn. Mater.* **189**, 214 (1998).
- <sup>34</sup> L. Li, K. Nishimura, W. D. Hutchison, Z. Qian, D. Huo, and T. Namiki, *Appl. Phys. Lett.* **100**, 152403 (2012).
- <sup>35</sup> Z. J. Mo, J. Shen, L. Q. Yan, J. F. Wu, L. C. Wang, J. Lin, C. C. Tang, and B. G. Shen, *Appl. Phys. Lett.* **102**, 192407 (2013).
- <sup>36</sup> W. Zuo, F. Hu, S. Jirong, and B. Shen, *J. Magn. Magn. Mater.* **344**, 96 (2013).
- <sup>37</sup> S. Gupta, R. Rawat, and K. G. Suresh, *Appl. Phys. Lett.* **105**, 012403 (2014).
- <sup>38</sup> H. Oesterreicher and F. T. Parker, *J. Appl. Phys.* **55**, 4334 (1984).
- <sup>39</sup> T. Krenke, E. Duman, M. Acet, E. F. Wassermann, X. Moya, L. Mañosa, and A. Planes, *Nat. Mater.* **4**, 450 (2005).