## Crystal structure and magnetic properties of R<sub>5</sub>Sn<sub>4</sub> alloys, where R is Tb, Dy, Ho, and Er

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Crystal structure and magnetic properties of  $R_5Sn_4$  alloys with R = Tb, Dy, Ho, and Er have been studied.  $R_5Sn_4$ ,  $R_{11}Sn_{10}$ , and  $R_5Sn_3$  phases coexist in the annealed alloys and the content of 11:10 and 5:3 phases varies between 9 and 17 wt.%. The  $R_5Sn_4$  major phase has  $Sm_5Ge_4$ -type orthorhombic structure with space group *Pnma*.  $Tb_5Sn_4$  has a complex magnetic structure, spin re-orientation and ferrimagnetic-paramagnetic phase transitions occur at ~54 and ~84 K, respectively. For Dy<sub>5</sub>Sn<sub>4</sub>, Ho<sub>5</sub>Sn<sub>4</sub>, and Er<sub>5</sub>Sn<sub>4</sub>, the antiferromagnetic-paramagnetic phase transitions occur at about 22, 15, and 8 K, respectively. The magnetic entropy changes ( $-\Delta S_M$ ) of all alloys are negative at low temperature and changes to positive at higher temperatures, which could be attributed to the change of magnetic states. © 2011 American Institute of Physics. [doi:10.1063/1.3549562]

The discovery of the giant magnetocaloric effect in Gd<sub>5</sub>Si<sub>2</sub>Ge<sub>2</sub> compound<sup>1</sup> has renewed the interest in rareearth/Si-group alloys<sup>2-5</sup> in search for materials suitable for magnetic refrigerants. However, information on the physics and chemistry of  $R_5T_4$  phases (T = Sn and Pb, and R = lanthanides except Gd) is still lacking.<sup>6</sup> Bulanova *et al.*<sup>7</sup> assessed the Ho-Sn binary phase diagram and showed existence of Ho<sub>5</sub>Sn<sub>4</sub> and Ho<sub>4</sub>Sn<sub>5</sub> compounds. They suggested that the Ho<sub>5</sub>Sn<sub>4</sub> compound could be formed by the peritectic reaction  $L + Ho_5Sn_3 \leftrightarrow Ho_5Sn_4$  at  $1720 \pm 13$  °C. Unfortunately, until now the existence of the compound Ho<sub>5</sub>Sn<sub>4</sub> has not been confirmed. Detecting Ho<sub>5</sub>Sn<sub>4</sub> phase by x-ray diffraction has not been successful because of the poor quality of the diffraction patterns in the intermediate range of compositions in the system.<sup>7</sup> As for Er-Sn phase diagram,<sup>8</sup> no Er<sub>5</sub>Sn<sub>4</sub> phase was reported; but a critical assessment from the available thermodynamic data of binary compounds of lanthanides and group IV elements suggested that both Ho<sub>5</sub>Sn<sub>4</sub> and Er<sub>5</sub>Sn<sub>4</sub> phases could exist.<sup>9</sup> This theoretical prediction calls for experimental verification. Here we report an experimental study of the crystal structure and magnetic properties of  $R_5Sn_4$  alloys with R = Tb, Dy, Ho, and Er.

The  $R_5Sn_4$  (R = Tb, Dy, Ho and Er) samples were prepared by arc-melting of constituent metals. 3 wt.% excess of Dy, Ho, and Er were added during sample preparation. The purity of Tb, Dy, Ho, Er and Sn was >99.99 wt.%. The ingots were arc melted several times to ensure homogeneity. The as-cast samples were annealed in a vacuum furnace at 1300 °C for 24 h. Powder x-ray diffraction (XRD) data were collected by using an X'Pert PRO diffractometer with Cu  $K\alpha_1$  radiation and analyzed by the Rietveld method using Rietica software.

The temperature (*T*) and magnetic field (*H*) dependencies of magnetization (*M*) were measured in a commercial superconducting quantum interference device magnetometer MPMS-XL (Quantum Design Inc). The ac magnetic susceptibility  $\chi_{ac}$  was measured with magnetic field amplitude of 2.5 Oe and frequency of 125 Hz. The isothermal magnetic entropy change,  $\Delta S_M$ , was calculated from the isothermal dc magnetization data using the Maxwell equation.<sup>10</sup>

The XRD refinement results of the phase contents in the as-cast and annealed  $R_5Sn_4$  (R = Tb, Dy, Ho, and Er) alloys are listed in Table I. The contents of 5:4 phases increase after annealing, especially in the Ho<sub>5</sub>Sn<sub>4</sub> and Er<sub>5</sub>Sn<sub>4</sub> alloys. The total amount of 11:10 and 5:3 phases varies between 9 and

TABLE I. The phase composition of the as-cast and annealed R5Sn4 alloys.

	Phase content (wt.%)			
Alloy	As-cast	Annealed-1300°/24 h		
Tb <sub>5</sub> Sn <sub>4</sub>	Tb <sub>5</sub> Sn <sub>4</sub> —82.3	Tb <sub>5</sub> Sn <sub>4</sub> —90.88		
	Tb <sub>11</sub> Sn <sub>10</sub> —11.97	Tb <sub>11</sub> Sn <sub>10</sub> -7.34		
	$Tb_5Sn_3$ —5.74	Tb <sub>5</sub> Sn <sub>3</sub> -1.78		
Dy <sub>5</sub> Sn <sub>4</sub>	Dy <sub>5</sub> Sn <sub>4</sub> —77.22	Dy <sub>5</sub> Sn <sub>4</sub> 82.22		
	Dy <sub>11</sub> Sn <sub>10</sub> —13.36	Dy <sub>11</sub> Sn <sub>10</sub> —10.28		
	Dy <sub>5</sub> Sn <sub>3</sub> —9.43	Dy <sub>5</sub> Sn <sub>3</sub> -7.50		
$Ho_5Sn_4$	Ho <sub>5</sub> Sn <sub>4</sub> —63.89	Ho <sub>5</sub> Sn <sub>4</sub> —83.54		
	Ho <sub>11</sub> Sn <sub>10</sub> —19.67	Ho <sub>11</sub> Sn <sub>10</sub> —7.10		
	Ho <sub>5</sub> Sn <sub>3</sub> —16.44	Ho <sub>5</sub> Sn <sub>3</sub> —9.36		
$Er_5Sn_4$	$Er_5Sn_4$ —64.93	Er <sub>5</sub> Sn <sub>4</sub>		
	$Er_{11}Sn_{10}$ —14.49	Er11Sn10-0.87		
	Er <sub>5</sub> Sn <sub>3</sub> —20.58	Er <sub>5</sub> Sn <sub>3</sub> —11.68		

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TABLE II.	Lattice parameters and	l unit cell	volumes of the R	2 <sub>5</sub> Sn <sub>4</sub> phases	with the orthorhombi	c Sm <sub>5</sub> Ge <sub>4</sub> -type structure.	

Alloy		Lattice parameters				
	<i>a</i> (Å)	<i>b</i> (Å)	<i>c</i> (Å)	$V(\text{\AA}^3)$	Space group	Remark
Tb <sub>5</sub> Sn <sub>4</sub>	8.004	15.4180	8.1390	1004.40	Pnma	Ref. [11]
	8.0086(3)	15.4369(7)	8.1390(4)	1006.2084(5)		This work
Dy <sub>5</sub> Sn <sub>4</sub>	7.9600	15.3600	8.1100	991.570	Pnma	Ref. [12]
	7.9865(5)	15.3693(9)	8.0930(5)	993.3749(1)		This work
Ho <sub>5</sub> Sn <sub>4</sub>	7.9628(2)	15.3001(4)	8.0534(2)	981.1631(8)	Pnma	This work
Er <sub>5</sub> Sn <sub>4</sub>	7.9366(1)	15.2330(3)	8.0140(2)	968.9095(4)	Pnma	This work



FIG. 1. (Color online) Temperature dependencies of magnetization of  $Tb_5Sn_4$  measured in the ZFC and FC modes (a). The insets of (a) show the temperature dependence of inverse dc magnetic susceptibility of  $Tb_5Sn_4$  measured in 100 Oe and 10 kOe applied fields. The magnetization in 100 Oe applied field of  $R_5Sn_4$  (R = Dy, Ho, and Er) alloys (b). The inset of (b) shows the temperature dependence of inverse dc magnetic susceptibility.

17 wt.%. The  $R_5Sn_4$  major phase in these alloys has the  $Sm_5Ge_4$ -type orthorhombic structure with space group *Pnma*. The lattice parameters and unit cell volumes of the  $R_5Sn_4$  phases are listed in Table II. The decrease of the lattice parameters and unit cell volumes with increasing rare earth atomic number reflects the lanthanide contraction.

Magnetic measurements were carried out on all annealed samples. The M(T) curve of Tb<sub>5</sub>Sn<sub>4</sub> alloy measured in 100 Oe and 10 kOe applied fields in "zero-field-cooled" (ZFC) and "field-cooled" (FC) modes are shown in Fig. 1(a). The insets display the temperature dependence of the inverse dc magnetic susceptibility (H/M). The Curie temperature  $T_{\rm C}$ , defined as the temperature at the maximum of |dM/dT| vs T plot, is  $\sim$ 84 K. The positive paramagnetic Curie temperature  $(\theta_p)$  obtained from the Curie Weiss fit and the magnetization isotherms shown in the inset of Fig. 2 indicate that the magnetic state of the Tb<sub>5</sub>Sn<sub>4</sub> alloy is ferrimagnetic. A cusp in the low field M(T) data and a peak in the high field data were observed at 54 K. Considering the amount of the major impurity phase ( $Tb_{11}Sn_{10}$ , 7.3 wt.%) and the large magnitude of the peak observed at 54 K with a 10 kOe field, it is unlikely that the cusp and peak are caused by the impurity phases. Similar to Tb<sub>5</sub>Ge<sub>4</sub>,<sup>13</sup> we assigned the cusp and peak observed at 54 K to a spin reorientation in Tb<sub>5</sub>Sn<sub>4</sub>. There is an



FIG. 2. (Color online) Temperature dependence of the ac magnetic susceptibility of  $R_5Sn_4$  (R = Tb, Dy, Ho, and Er) alloys. The inset shows magnetization isotherms of  $Tb_5Sn_4$  alloy with increasing magnetic field.



FIG. 3. (Color online) Temperature dependence of isothermal magnetic entropy changes  $(-\Delta S_M)$  of  $R_5Sn_4$  (R = Tb, Dy, Ho, and Er) alloys for various magnetic field changes.

additional kink near 16 K [Fig. 1(a)], whose origin needs further investigation.

The temperature dependencies of dc magnetization measured with a 100 Oe field for the  $R_5Sn_4$  (R = Dy, Ho, and Er) alloys are shown in Fig. 1(b). The M(T) curves exhibit a typical antiferromagnetic behavior, which agrees with the negative  $\theta_p$  values as shown in the inset of Fig. 1(b). The Néel temperatures are 22, 14, and 8 K for Dy<sub>5</sub>Sn<sub>4</sub>, Ho<sub>5</sub>Sn<sub>4</sub>, and Er<sub>5</sub>Sn<sub>4</sub> alloys, respectively.

The temperature dependencies of the ac magnetic susceptibility for  $R_5Sn_4$  (R = Tb, Dy, Ho, and Er) alloys are displayed in Fig. 2. For  $Tb_5Sn_4$  alloy, a distinct anomaly (peak) is observed at  $T_C$  (~84 K), and a cusp near ~54 K indicates a spin-reorientation. The results are in good agreement with the temperature dependent dc magnetization measured with a 100 Oe applied field. The antiferromagnetic-paramagnetic phase transitions occur at about 22, 15, and 8 K for Dy<sub>5</sub>Sn<sub>4</sub>, Ho<sub>5</sub>Sn<sub>4</sub>, and Er<sub>5</sub>Sn<sub>4</sub> alloys, respectively. The Néel points determined from the ac and dc magnetic susceptibility are in good agreement.

Figures 3(a)–3(d) show the  $(-\Delta S_M) \sim T$  for different magnetic field changes for Tb<sub>5</sub>Sn<sub>4</sub>, Dy<sub>5</sub>Sn<sub>4</sub>, Ho<sub>5</sub>Sn<sub>4</sub>, and Er<sub>5</sub>Sn<sub>4</sub>, respectively. The  $-\Delta S_M$  vs *T* plot of Tb<sub>5</sub>Sn<sub>4</sub> alloy exhibits a negative plateau between 20 and 52 K. It changes sign and exhibits a positive plateau at higher temperatures (60 ~ 95 K). The sign change occurs in the vicinity of the spin re-orientation transition. Figs. 3(b)–3(d) indicate that the magnetic entropy changes of Dy<sub>5</sub>Sn<sub>4</sub>, Ho<sub>5</sub>Sn<sub>4</sub>, and Er<sub>5</sub>Sn<sub>4</sub> show a similar negative to positive sign change behavior. Since the ( $-\Delta S_M$ ) of an antiferromagnet is negative; while for a ferromagnet, it is positive.<sup>14</sup> The change of the sign of the  $-\Delta S_M$  in our R<sub>5</sub>Sn<sub>4</sub> alloys is likely due to a change of the magnetic structure between antiferromagnetically and ferromagnetically dominated states caused by the applied field. The maximum  $-\Delta S_M$  values for these  $R_5Sn_4$  alloys with R = Tb, Dy, Ho, and Er are 1.49, 2.92, 7.95, and 10.25 J/kg K, respectively, with a field change of 50 kOe.

We note that the  $Tb_5Sn_4$  alloy exhibits thermal irreversibility with low applied field, spin reorientation, and ferrimagnetic behaviors, which are different from  $Dy_5Sn_4$ ,  $Ho_5Sn_4$ , and  $Er_5Sn_4$  alloys. Noticeable differences in the magnetic properties among  $R_5Ge_4$  (R = Tb, Dy, Ho, and Er) were also observed. Such differences are not surprising given the differences of these rare earth ions. To understand underlying mechanism, detailed magnetic structure information on these alloys is needed.

In summary, 5:4, 11:10, and 5:3 phases coexist in the annealed  $R_5Sn_4$  alloys with R = Tb, Dy, Ho, and Er. The major  $R_5Sn_4$  phase in these alloys has the  $Sm_5Ge_4$ -type orthorhombic structure with space group *Pnma*. The content of 11:10 and 5:3 phases varies between 9 and 17 wt.%. Both the dc magnetization and ac susceptibility indicate that  $Tb_5Sn_4$  alloy has a complex magnetic structure. The spin reorientation and ferrimagnetic-paramagnetic phase transitions occur near 54 and 84 K, respectively. For Dy<sub>5</sub>Sn<sub>4</sub>, Ho<sub>5</sub>Sn<sub>4</sub>, and Er<sub>5</sub>Sn<sub>4</sub> alloys, the antiferromagnetic-paramagnetic phase transition occurs at about 22, 15, and 8 K respectively. The  $(-\Delta S_M)$  of all alloys change sign from negative to positive with increasing temperature, which could be attributed to the change of the magnetic state of the compound.

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- <sup>1</sup>V. K. Pecharsky and K. A. Gschneidner, Jr., Phys. Rev. Lett. 78, 4494 (1997).
- <sup>2</sup>L. Morellon, C. Magen, P. A. Algarabel, M. R. Ibarra, and C. Ritter, Appl. Phys. Lett. **79**, 1318 (2001).
- <sup>3</sup>V. V. Ivtchenko, V. K. Pecharsky, and K. A. Gschneidner, Jr., Adv. Cryog. Eng. Mater. **46**, 405 (2000).
- <sup>4</sup>A. O. Pecharsky, K. A. Gschneidner, Jr., V. K. Pecharsky, D. L. Schlagel, and T. A. Lograsso, Phys. Rev. B 70, 144419 (2004).
- <sup>5</sup>D. H. Ryan, M. Elouneg-Jamróz, J. van Lierop, Z. Altounian, and H.B. Wang, Phys. Rev. Lett. **90**, 117202 (2003).
- <sup>6</sup>V. K. Pecharsky and K. A. Gschneidner, Jr., Pure Appl. Chem. **79**, 1383 (2007).
- <sup>7</sup>M. V. Bulanova, V. N. Eremenko, V. M. Petjukh, and V. R. Sidorko, J. Phase Equilib. **19**, 136 (1998).
- <sup>8</sup>H. Okamoto, J. Phase Equilib. 16, 197 (1995).
- <sup>9</sup>V. T. Witusiewicz, V. R. Sidorko, and M. V. Bulanova, J. Alloys Compd. **248**, 233 (1997).
- <sup>10</sup>V. K. Pecharsky and K. A. Gschneidner, Jr., J. Appl. Phys. 86, 565 (1999).
- <sup>11</sup>V. N. Eremenko, M. V. Bulanova, and P. S. Martsenyuk, Sov. Prog. Chem. **57**, 14 (1990). Calculated from LPF using POWD-12++.
- <sup>12</sup>V. N. Eremenko, M. V. Bulanova, and P. S. Martsenyuk. J. Alloys Compds. 189, 229 (1992). Calculated from LPF using POWD-12++.
- <sup>13</sup>C. Ritter, L. Morellon, P. A. Algarabel, C. Magen, and M. R. Ibarra, Phys. Rev. B 65, 094405 (2002).
- <sup>14</sup>A. M. Tishin and Y. I. Spichkin, The Magnetocaloric Effect and Its Applications (Institute of Physics, Bristol, 2003).